species B

 $\alpha, \alpha =$ lower and upper solutions of α

 $\overline{\beta}$, $\overline{\beta}$ = lower and upper solutions of β θ = dimensionless variable of time ξ = dimensionless variable of the space.

= dimensionless variable of the spatial coordinate x

= ratio of the coefficient of diffusion of species A

to the coefficient of diffusion of species B

= ratio of the concentration of species B at t = 0г to the concentration of species \bar{A} at x = 0

 G_p , $R_p = \text{domain and boundary of systems (1) and (2)}$

 R_n = subset of R_p R_p^+ , R_* = boundary, defined in definition 1

 θ_k , ξ_k = sequences of θ and ξ , defined in definition 1

 $\phi, \psi = \text{functions defined in } G_p, \text{ see definition 1 } ||\xi_k||_e = \text{Euclidean distance}$

f, g = functions, defined by (7) and (8) D(f), D(g) = domains of f and g $\rho = \text{function}$, defined in definition 2 $\tau, \delta = \text{small number}$, defined in definition 2

T = some large dimensionless variable of time $Z_0(f)$, $Z_0(g, R_n)$ = function class, defined in definitions

= function, defined by (19)

 P_1 , P_2 = differential operators, defined in (7) and (8)

U = union, set notation

= element of, set notation = subset of, set notation

= function class, defined in definition 2

LITERATURE CITED

Brian, P. L. T., J. F. Hurley, and E. H. Hasseltine, "Penetration Theory for Gas Absorption Accompanied by a Second Order Chemical Reaction," AIChE J., 7, 226 (1962). Carslaw and Jaeger, Conduction of Heats in Solids, pp. 63,

134, Clarendon Press, Oxford, England (1959).

McDaniel, T. J., and V. R. Murthy, "Solution Bounds to Structural Systems," AIAA J., 14, 111 (1976).

—————, "Bounds on the Dynamic Characteristics of Ro-

tating Beams," ibid., 439 (1977).

Pearson, J. R. A., "Diffusion of one Substance into a Semi-Infinite Medium Containing Another with Second Order Reaction." Appl. Sci. Res., A11, 321 (1963).

Tam, K. K., "On the Oseen Linearization of the Swirling Flow

Boundary Layer," J. Fluid Mech., 75, 777 (1976).

—, and K. Y. K. Ng, "Construction of Upper and Lower Solutions for Flow Past a Non-Uniformly Heated Plate," J. Math. Anal. Application, 59, 531 (1977).

Walter, W., Differential and Integral Inequalities, Springer Verlag, New York (1970).

Manuscript received February 9, 1978; revision received June 5, and accepted June 20, 1978.

Prediction of Partial Molar Volume from the Lee-Kesler Equation of State

C.T. LIN and T. E. DAUBERT

Department of Chemical Engineering The Pennsylvania State University University Park, Pennsylvania 16802

$$D = d_1 + \frac{d_2}{T_r} \tag{5}$$

and b_1 , b_2 , b_3 , b_4 , c_1 , c_2 , c_3 , c_4 , d_1 , d_2 , β , and γ are constants. The set of constants for the simple fluid are determined from the data for argon, krypton, and methane and that of the reference fluid are from the data for n-octane. The Lee-Kesler correlation is accurate in estimating thermodynamic properties for nonpolar fluids (Reid et al., 1977; Danner et al., 1976). In this work, the Lee-Kesler equation of state is applied in the calculation of the partial molar volumes of nonpolar binary mixtures.

(1)

The function for both the simple fluid $F^{(0)}$ and the reference fluid $F^{(r)}$ are derived from a reduced form of the modified BWR equation of state with a different

Based on Pitzer's three-parameter corresponding states

principles (1955), Lee and Kesler (1975) developed an analytical method to represent the thermodynamic function of any fluid F in terms of a simple fluid contribution

 $F^{(o)}$ and of a reference fluid contribution $F^{(r)}$ as follows:

$$Z = \left(\frac{P_r V_r}{T_r}\right) = 1 + \frac{B}{V_r} + \frac{C}{V_r^2} + \frac{D}{V_r^3} + \frac{C_4}{T_r^3 V_r^3} \left(\beta + \frac{\gamma}{V_r^2}\right) \exp\left(-\frac{\gamma}{V_r^2}\right) \quad (2)$$

where

$$B = b_1 - \frac{b_2}{T_r} - \frac{b_3}{T_{r^2}} - \frac{b_4}{T_{r^3}}$$
 (3)

$$C = c_1 - \frac{c_2}{T_r} + \frac{c_3}{T_r^3} \tag{4}$$

0001-1541-79-1670-0365-\$00.75. © The American Institute of Chemical Engineers, 1979.

METHOD OF CALCULATION

If the thermodynamic function F in Equation (1) is the compressibility factor, then Equation (1) can be written as

$$Z = Z^{(o)} + \frac{\omega}{\omega^{(r)}} [Z^{(r)} - Z^{(o)}]$$
 (6)

From the definition, $Z = P_{\nu}/RT$, $Z^{(o)} = P_{r}V_{r}^{(o)}/T_{r}$, and $Z^{(r)} = P_r V_r^{(r)} / T_r$, we obtain

$$\nu = \frac{RT}{P} \left\{ \frac{P_r}{T_r} \left[V_r^{(o)} + \frac{\omega}{\omega^{(r)}} \left(V_r^{(r)} - V_r^{(o)} \right) \right] \right\}$$
(7)

Table 1. Comparison of the Partial Molar Volume Estimated by the Lee-Kesler Equation of State and That Calculated from Data on Volume Change

Partial molar volume, cm3/g-mole

	Tempera-	Pressure,	Composi-	Component 1		Component 2		
System	ture, °K	atm	tion, x_1	Estimate	Experiment	Estimate	Experiment	Data and source
Methane (1) n-butane (2) Ethane (1)	344.3	40	0.1487	93.31	101.48	111.39	107.9	Reamer et al (1947)
n-pentane (2) Benzene (1)	344.3	272	0.6250	74.58	78.10	115.49	101.58	Sage (1965)
2,2,4, trimethyl- pentane (2) Benzene (1)	293.33	1	0.5	89.42	88.87	160.81	166.19	Graboski (1970)
1-octene (2)	293.33	1	0.5	89.91	89.02	155.33	157.70	Graboski (1970)
Benzene (1) Ethylcyclohexane (2) Benzene (1)	293.33	1	0.5	89.72	89.06	150.95	143.48	Graboski (1970)
Benzene (1) n-dodecane (2)	293.33	1	0.5	87.21	89.49	223.65	228.71	Graboski (1970)

The total volume V is equal to the total number of moles multiplied by the molar volume:

If derivatives of Equation (10) with respect to N_i are taken, and the final equation is rearranged

$$\frac{\left(\frac{\partial T_r}{\partial N_i}\right) \left[-\frac{1}{V_r} \left(-\frac{b_2}{T_r^2} + \frac{2b_3}{T_r^3} + \frac{3b_4}{T_r^4} \right) + \frac{1}{V_r^2} \left(-\frac{c_2}{T_r^2} - \frac{3c_3}{T_r^4} \right) - \frac{d_2}{V_r^5 T_r^2} + \frac{P_r V_r}{T_r^2} - \frac{3c_4}{T_r^4 V_r^2} \right] }{\left\{ -\frac{P_r}{T_r} + \frac{B}{V_r^2} + \frac{2C}{V_r^3} + \frac{5D}{V_r^6} + \exp\left(-\frac{\gamma}{V_r^2} \right) \left[-\frac{2c_4}{T_r^3 V_r^3} \left(\beta + \frac{\gamma}{V_r^2} \right) + \frac{2\gamma c_4}{T_r^3 V_r^5} \right] \right\} }$$

$$V = \frac{N_T R I'}{P} \left\{ \frac{P_r}{T_r} \left[V_r^{(o)} + \frac{\omega}{\omega^{(r)}} \left(V_r^{(r)} - V_r^{(o)} \right) \right] \right\}$$
(8)

By the definition, the partial molar volume of component i in a mixture can be derived from Equation (8):

$$\overline{V}_{i} = \left(\frac{\partial V}{\partial N_{i}}\right)_{T,P,N_{i}(j \neq i)} = \frac{ZRT}{P} \left\{ 1 - \frac{N_{T}}{P_{c}} \left(\frac{\partial P_{c}}{\partial N_{i}}\right) \right\}$$

Since

$$\left(\frac{\partial T_r}{\partial N_i}\right) = -\frac{T}{T_c^2} \left(\frac{\partial T_c}{\partial N_i}\right) \tag{12}$$

$$\frac{\partial P_r}{\partial N_i} = -\frac{P}{P_c^2} \left(\frac{\partial P_c}{\partial N_i} \right) \tag{13}$$

Substituting Equations (12) and (13) into Equation (11), we have

$$(\partial V_r/\partial N_i) = \frac{\frac{P_r V_r}{T_r} \frac{1}{P_c} \left(\frac{\partial P_c}{\partial N_i}\right) - T_r \frac{1}{T_c} \left(\frac{\partial T_c}{\partial N_i}\right) \left[\frac{1}{V_r} \left(\frac{b_2}{T_r^2} + \frac{2b_3}{T_r^3} + \frac{3b_4}{T_r^4}\right) + \frac{1}{V_r^2} \left(\frac{c_2}{T_r^2} - \frac{3c_3}{T_r^4}\right) + \frac{P_r V_r}{T_r} \right]}{\left\{\frac{P_r}{T_r} + \frac{B}{V_r^2} + \frac{2C}{V_r^3} + \frac{5D}{V_r^6} + \exp\left(-\frac{\gamma}{V_r^2}\right) \left[\frac{2c_4}{T_r^3 V_r^3} \left(\beta + \frac{\gamma}{V_r^2}\right) + \frac{2\gamma c_4}{T_r^3 V_r^5}\right]\right\}}$$

$$+ \frac{N_{T}}{T_{c}} \left(\frac{\partial T_{c}}{\partial N_{i}} \right) + \frac{N_{T}RT_{c}}{P_{c}} \left\{ \left(\frac{\partial V_{r}^{(o)}}{\partial N_{i}} \right) + \frac{\omega}{\omega^{(r)}} \left[\left(\frac{\partial V_{r}^{(l)}}{\partial N_{i}} \right) - \left(\frac{\partial V_{r}^{(o)}}{\partial N_{i}} \right) \right] + \frac{1}{\omega^{(r)}} \left(V_{r}^{(l)} - V_{r}^{(o)} \right) \left(\frac{\partial \omega}{\partial N_{i}} \right) \right\}$$
(9)

In order to calculate the derivatives of reduced volume of both the simple and the reference fluid, $(\partial V_r^o/\partial N_i)$ and $(\partial V_r^1/\partial N_i)$, respectively, Equation (2) is rearranged as

$$P_{r}V_{r} = T_{r} \left\{ 1 + \left(b_{1} - \frac{b_{2}}{T_{r}} - \frac{b_{3}}{T_{r}^{2}} - \frac{b_{4}}{T_{r}^{3}} \right) \frac{1}{V_{r}} + \left(c_{1} - \frac{c_{2}}{T_{r}} + \frac{c_{3}}{T_{r}^{3}} \right) \frac{1}{V_{r}^{2}} + \left(d_{1} + \frac{d_{2}}{T_{r}} \right) \frac{1}{V_{r}^{5}} + \frac{c_{4}}{T_{r}^{3}V_{r}^{2}} \left(\beta + \frac{\gamma}{V_{r}^{2}} \right) \exp\left(-\frac{\gamma}{V_{r}^{2}} \right) \right\}$$
(10)

Equations for calculating $(\partial P_c/\partial N_i)$, $(\partial T_c/\partial N_i)$, $(\partial \omega/\partial N_i)$, and $(\partial V_c/\partial N_i)$ are given in the Appendix. If constants of simple fluids and reference fluids for Equation (2) are substituted into Equation (13), $(\partial V_r^{(o)}/\partial N_i)$ and $(\partial V_r^{(1)}/\partial N_i)$ are obtained. Equation (9) can then be utilized to estimate the partial molar volumes.

RESULTS

Table 1 compares the partial molar volumes estimated by the Lee-Kesler equation with experimental values. The estimates of Equation (9) are very close to those calculated from the data of volume change due to the variation of concentration in the mixture. Even at a pressure of 272 atm, the accuracy of the predicted values are still within 15% of the experimental values. Generally speaking, Equation (9) can predict accurate partial molar volumes of a component in a nonpolar binary mixture at moderate temperature and pressure. However, when the mixture approaches within about 5% of a pure component, the accuracy of the estimates of Equation (9) cannot be verified. Either the data for the volume change are not available, or they are subject to uncertainty.

APPENDIX

Equations to calculate the derivatives of the critical properties with respect to molar number.

For a binary system which contains N_1 moles of component 1 and N_2 moles of component 2, the critical volume of the mixture is

$$V_c = \frac{1}{N_T^2} \left\{ N_1^2 V_{c1} + N_2^2 V_{c2} + \frac{1}{4} N_1 N_2 (V_{c1}^{1/3} + V_{c2}^{1/3})^3 \right\}$$
(A1)

where $N_{\rm T} = N_1 + N_2$.

Taking derivatives with respect to N_1 and N_2 , respectively, the following two equations result:

$$\left(\frac{\partial V_c}{\partial N_1}\right) = \frac{1}{N_T} \left\{ 2(x_1 V_{c_1} - V_c) + \frac{1}{4} x_2 (V_{c_1}^{1/3} + V_{c_2}^{1/3})^3 \right\}$$
(A2)

$$\left(\frac{\partial V_c}{\partial N_2}\right) = \frac{1}{N_T} \left\{ 2(x_2 V_{c_2} - V_c) + \frac{1}{4} x_1 (V_{c_1}^{1/3} + V_{c_2}^{1/3})^3 \right\}$$
(A3)

Here $x_1 \equiv N_1/N_T$ and $x_2 \equiv N_2/N_T$.

$$\frac{\left(\beta + \frac{\gamma}{V_r^2}\right) \exp\left(-\frac{\gamma}{V_r^2}\right) \left] - \frac{V_r}{T_r} \left(\frac{\partial P_r}{\partial N_i}\right)}{\left(1 - \beta - \frac{\gamma}{V_r^2}\right)\right]} \tag{11}$$

According to the mixing rules recommended by Lee and Kesler, the critical temperature of the mixture is

$$T_c = -\frac{1}{N_T^2 V_c} \left\{ N_1^2 V_{c_1} T_{c_1} + \frac{1}{4} N_1 N_2 (V_{c_1}^{1/3} + V_{c_2}^{1/3})^3 \right.$$

$$\sqrt{T_{c_1} T_{c_2}} + N_2^2 V_{c_2} T_{c_2} \left. \right\}$$
 (A4)

$$\frac{-\frac{3c_4}{T_r^4 V_r^2} \left(\beta + \frac{\gamma}{V_r^2}\right) \exp\left(-\frac{\gamma}{V_r^2}\right) \right]}{\left(1 - \beta - \frac{\gamma}{V_r^2}\right) \right]}$$
(14)

The following two equations are obtained by taking derivatives with respect to N_1 and N_2 , respectively:

$$\left(\frac{\partial T_c}{\partial N_1}\right) = \frac{1}{N_T V_c} \left\{ 2(x_1 V_{c_1} T_{c_1} - V_c T_c) + \frac{1}{4} x_2 (V_{c_1}^{1/3} + V_{c_2}^{1/3})^3 \sqrt{T_{c_1} T_{c_2}} - T_c N_T \left(\frac{\partial V_c}{\partial N_1}\right) \right\} (A5)$$

$$\left(\frac{\partial T_c}{\partial N_2}\right) = \frac{1}{N_T V_c} \left\{ 2(x_2 V_{c_2} T_{c_2} - V_c T_c) + \frac{1}{4} x_1 (V_{c_1}^{1/3} + V_{c_2}^{1/3})^3 \sqrt{T_{c_1} T_{c_2}} - T_c N_T \left(\frac{\partial V_c}{\partial N_2}\right) \right\} (A6)$$

Since the critical pressure of the mixture is

$$P_c = (0.2905 - 0.085 \,\omega) \,\frac{RT_c}{V_c} \tag{A7}$$

taking derivatives of Equation (A7) with respect to N1 and

$$\left(\frac{\partial P_c}{\partial N_1}\right) = -0.085 \left(\frac{\partial \omega}{\partial N_1}\right) \frac{RT_c}{V_c} + (0.2905 - 0.085 \,\omega)^*$$

$$\left[\left(\frac{\partial T_c}{\partial N_1}\right) - \frac{T_c}{V_c}\left(\frac{\partial V_c}{\partial N_1}\right)\right] \frac{R}{V_c} \quad (A8)$$

$$\left(\frac{\partial P_c}{\partial N_2}\right) = -0.085 \left(\frac{\partial \omega}{\partial N_2}\right) \frac{RT_c}{V_c} + (0.2905 - 0.085 \,\omega)^{\bullet}$$

$$\left[\left(\frac{\partial T_c}{\partial N_2} \right) - \frac{T_c}{V_c} \left(\frac{\partial V_c}{\partial N_2} \right) \right] \frac{R}{V_c} \quad (A9)$$

where R is the gas constant

$$\left(\frac{\partial \omega}{\partial N_1}\right) = \frac{1}{N_T} \left\{\omega_1 - \omega\right\} \tag{A10}$$

$$\left(\frac{\partial \omega}{\partial N_2}\right) = \frac{1}{N_T} \left\{\omega_2 - \omega\right\} \tag{A11}$$

According to the mixing rules proposed by Lee and Kesler, the acentric factor of the mixture is

$$\omega = \frac{1}{N_T} \left\{ N_1 \omega_1 + N_2 \omega_2 \right\} \tag{A12}$$

NOTATION

molar number of component i in the mixture

total molar number in the system

pressure

critical pressure of mixture

critical pressure of component i

 P_r T T_c T_{ci} T_r V= reduced pressure = temperature

= critical temperature of mixture

= critical temperature of component i

= reduced temperature

= total volume

= critical volume of the mixture = critical volume of component i

= reduced volume = compressibility

= molar volume

acentric factor

accentric factor of reactor fluid

LITERATURE CITED

Danner, R. P., T. E. Daubert, T. G. Katinas, C. A. Passut, and R. R. Tarakad, Documentation of the Basis for Selection of the Contents of Chapter 7 of API Technical Data Book-

Petroleum Refining, Pa. State Univ., University Park (1976). Lee, B. I., and M. G. Kesler, "A Generalized Thermodynamic Correlation Based on Three-Parameter Corresponding States," AIChE J., 21, 510 (1975). Graboski, M. S., "Physical Properties of Binary Hydrocarbon

Mixtures," M.S. thesis, Pa. State Univ., University Park (1970).

Pitzer, K. S., D. Z. Lippman, R. F. Curl, Jr., C. M. Huggins, and D. E. Peterson, "The Volumetric and Thermodynamic Properties of Fluids. II. Compressibility Factor, Vapor Pressure and Entropy of Vaporization," J. Am. Chem. Soc., 77, 3433 (1955).

Reamer, H. H., K. J. Korpi, B. H. Sage, and W. N. Lacey, "Phase Equilibria in Hydrocarbon Systems. Volumetric Behavior of Methane-n-Butane Systems at Higher Pressures," Ind. Eng. Chem., 39, 206 (1947).

Reid, R. C., J. M. Prausnitz, and T. K. Sherwood, The Properties of Gases and Liquids, 3 ed., McGraw-Hill, New York (1977)

Sage, B. H., Thermodynamics of Multicomponent Systems, pp. ew York, Chapman & Hall, Ltd., London, England (1965).

Manuscript received April 6, 1978; revision received June 9, and accepted June 20, 1978.